## ON BOUNDARIES OF THE GROWTH REGION OF MULTIDIMENSIONAL PERTURBATIONS OF UNSTABLE STATES

PMM Vol.41, № 6, 1977, pp. 1122-1123

V. A. KONDRASHEV and A. G. KULIKOVSKII
(Moscow)
(Received December 14, 1976)
When the development of instability in the coordinate-independent stationary states is investigated, the determination of the boundary of the expanding region within which the perturbations grow, initially specified in a bounded region is of great importance. In particular, the knowledge of these boundaries makes it possible to decide whether the instability is absolute or convective [1, 2]. Below it will be shown that a boundary of the region occupied by growing perturbations can be obtained, in the non-onedimensional case, in the form of an envelope of the straight lines or planes bounding the region in which the growth of the one-dimensional perturbations takes place.
We shall limit ourselves, for convenience, to investigating the two-dimensional perturbations (the three-dimensional perturbations are investigated in the same manner). As we know, a perturbation which is initially localized and represented by the Fourier integral [3]

$$
u(x, y, t)=\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \exp \left[i k_{1} x+i k_{2} y-i \omega\left(k_{1}, k_{2}\right) t\right] d k_{1} d k_{2}
$$

can be estimated along the rays $x=U t, y=V t \quad$ as $\quad t \rightarrow \infty$, using the method of steepest descent [3] in accordance with the formulas

$$
\begin{align*}
& \exp \left[t \operatorname{Im} \omega^{\prime}(U, V)-\ln t\right]  \tag{1}\\
& \omega^{\prime}=\omega\left(k_{1}, k_{2}\right)-k_{1} U-k_{2} V  \tag{2}\\
& \partial \omega / \partial k_{1}=U, \quad \partial \omega / \partial k_{2}=V \tag{3}
\end{align*}
$$

where $\omega=\omega\left(k_{1}, k_{2}\right)$ represents the dispersion equation of the perturbations. We find the function $\omega^{\prime}=\omega^{\prime}(U, V)$ by obtaining the values of $k_{1}$ and $k_{2}$ from (3), the latter defining the points of steepest descent of the function: $\omega^{\prime}\left(k_{1}, k_{2}, U, V\right)$ on the complex planes $k_{1}$ and $k_{2}$, and substituting these values into (2).

The curve defined by the equation

$$
\begin{equation*}
\operatorname{lm} \omega^{\prime}(U, V)=0 \tag{4}
\end{equation*}
$$

separates the domain of values of $U$ and $V$ at which the perturbations grow along the ray $x=U t, y=V t$, from the domain corresponding to the decay of the perturbations.

Let us consider the one-dimensional perturbations corresponding to the same dispersion equation and such, that a coordinate system can be chosen so that $\operatorname{Im} k_{2}=0$.

The perturbations in this coordinate system will not grow along the $y$-axis, and we shall regard the real quantity $k_{2}$ as a parameter. The asymptotic behavior of the one-dimensional perturbations along the ray $x=U t$ as $t \rightarrow \infty$ will have the form

$$
\begin{align*}
& \exp \left[t \operatorname{Im} \omega^{\prime \prime}\left(k_{2}, U\right)-1 / 2 \ln t\right]  \tag{5}\\
& \omega^{\prime \prime}-\omega\left(k_{1}, k_{2}\right)-k_{1} U, \quad \partial \omega / \partial k_{1}=U \tag{6}
\end{align*}
$$

The boundaries of the domain on the $U$-axis corresponding to the growing perturbations are determined by the values of $U$ extremal in $k_{2}$ and satisfying the equation

$$
\begin{equation*}
\operatorname{Im} \omega^{\prime \prime}\left(k_{2}, U\right)=0 \tag{7}
\end{equation*}
$$

Let the quantity $U$ assume, at certain $k_{2}=k_{20}$, its extremal value $U_{0}$. We shall show that a point can be found on the straight line $U=U_{0}$ lying in the $U, V$-plane, which belongs to the boundary curve (4). When $k_{2}=k_{20}$ and $U=U_{0}$, we have

$$
\begin{equation*}
\left(\partial \operatorname{Im} \omega^{\prime \prime} / \partial k_{2}\right)_{U}=0 \tag{8}
\end{equation*}
$$

According to (6) we have

$$
\begin{equation*}
\left(\frac{\partial \omega^{\prime \prime}}{\partial k_{2}}\right)_{V}=\left(\frac{\partial \omega}{\partial k_{1}}\right)\left(\frac{\partial k_{1}}{\partial k_{2}}\right)_{V}+\left(\frac{\partial \omega}{\partial k_{2}}\right)_{k_{1}}-U\left(\frac{\partial k_{1}}{\partial k_{2}}\right)=\left(\frac{\partial \omega}{\partial k_{2}}\right)_{k_{1}} \tag{9}
\end{equation*}
$$

From (8) and (9) it follows that a $V_{0}$ exists such that

$$
\begin{equation*}
\left(\partial \omega / \partial h_{2}\right)_{h_{1}}=V_{0}, \quad \operatorname{Im} V_{0}=0 \tag{10}
\end{equation*}
$$

The fact that $k_{2}$ and $V_{0}$ are real implies, together with (7), that the relation (4) holds for the values $U_{0}$ and $V_{0}$, i. e. that the latter point belongs to the boundary curve.

We shall now show that the straight line $U=U_{\mathrm{n}}$ touches the boundary curve (4) at the point $U_{0}, V_{0}$. Indeed, the derivative of $\operatorname{Im} \omega^{\prime}$ along the direction of
$V$ will vanish only under the condition that the curve (4) has a vertical tangent. Using the relations (2) and (3), we obtain

$$
\begin{aligned}
& \left(\frac{\partial \operatorname{Im} \omega^{\prime}}{\partial V}\right)_{\Gamma}=\operatorname{lm}\left[\left(\frac{\partial \omega}{\partial h_{1}}\right)_{R_{2}}\left(\frac{\partial h_{1}}{\partial V}\right)_{I} \therefore\left(\frac{\partial \omega}{\partial h_{2}}\right)_{k_{1}}\left(\frac{\partial h_{2}}{\partial V}\right)_{r}-k_{2} \cdots U \times\right. \\
& \left.\left(\frac{\partial k_{1}}{\partial V^{\prime}}\right)_{\Gamma}-V\left(\frac{\partial h_{2}}{\partial V}\right) \Gamma \right\rvert\,=-\operatorname{lm} k_{2}
\end{aligned}
$$

Considering now the straight lines of all possible orientations bounding the regions of growth of the one-dimensional perturbations, we obtain (4) as their envelope.

REFERENCES

1. Dysthe, K. B., Convective and absolute instability. Nucl. Fusion, Vol. 6, No. 3, 1966.
2. Fainberg, Ia. B., Kurilko, V.I. and Shapiro, V. D., On the problem of insta bilities in the course of interaction of charged particle beams with plasma. Zh. tekhn. fiz. Vol. 31, No. 6, 1961.
3. Evgrafov, M. A. and Postnikov, M. M., Asymptotics of the Green's functions of parabolic and elliptic equations with constant coefficients. Matem. sb. Vol. 82, No. 1, 1970. Translated by L. K.
